

RL Circuits

"Building-Up" Phase:

Connecting the switch to **position A** corresponds to the **"building up" phase of an RL circuit**. Summing all the potential changes in going around the loop gives

$$\varepsilon - IR - L \frac{dI}{dt} = 0 ,$$

where **I(t)** is a function of time. If the switch is closed (**position A**) at $t=0$ and **I(0)=0** (assuming the current is zero at $t=0$) then

$$\frac{dI}{dt} = - \frac{1}{\tau} \left(I - \frac{\varepsilon}{R} \right) , \quad \text{where I have define } \tau=L/R.$$

Dividing by $(I-\varepsilon/R)$ and multiplying by dt and integrating gives

$$\int_0^I \frac{dI}{(I - \varepsilon / R)} = - \int_0^t \frac{1}{\tau} dt , \quad \text{which implies } \ln \left(\frac{I - \varepsilon / R}{-\varepsilon / R} \right) = - \frac{t}{\tau} .$$

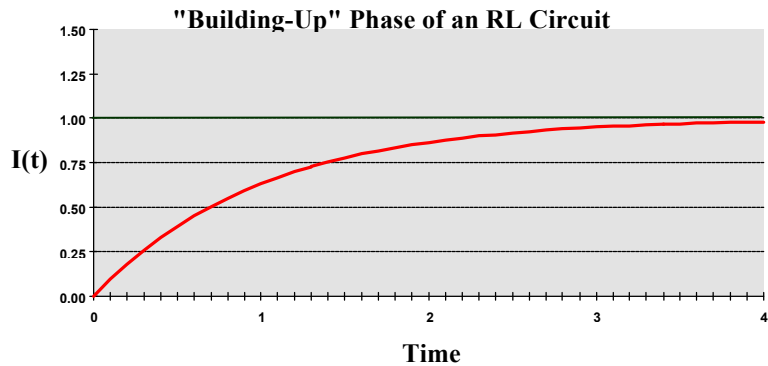
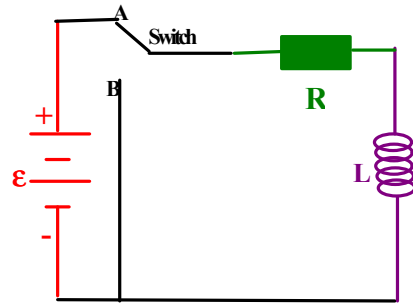
Solving for **I(t)** gives

$$I(t) = \frac{\varepsilon}{R} (1 - e^{-t/\tau}) .$$

The potential change across the inductor is given by $\Delta V_L(t)=-LdI/dt$ which yields

$$\Delta V_L(t) = -\varepsilon e^{-t/\tau} .$$

The quantity $\tau=L/R$ is call the **time constant** and has dimensions of time.



"Collapsing" Phase:

Connecting the switch to **position B** corresponds to the **"collapsing" phase of an RL circuit**. Summing all the potential changes in going around the

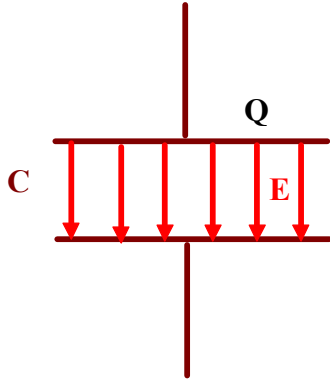
loop gives $- IR - L \frac{dI}{dt} = 0$, where **I(t)** is a function of time. If the

switch is closed (**position B**) at $t=0$ then **I(0)=I₀** and

$$\frac{dI}{dt} = - \frac{1}{\tau} I \quad \text{and} \quad I(t) = I_0 e^{-t/\tau} .$$

Capacitors and Inductors

Capacitors Store Electric Potential Energy:

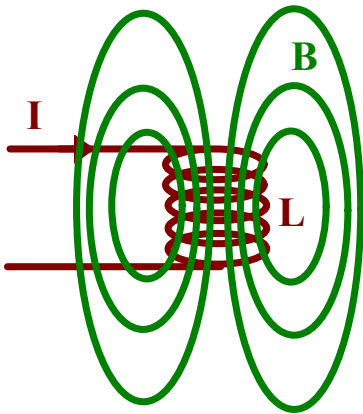


$$U_E = \frac{Q^2}{2C}$$

$$Q = C\Delta V_C \quad \Delta V_C = Q / C$$

$$u_E = \frac{1}{2} \epsilon_0 E^2 \quad (\text{E-field energy density})$$

Inductors Store Magnetic Potential Energy:



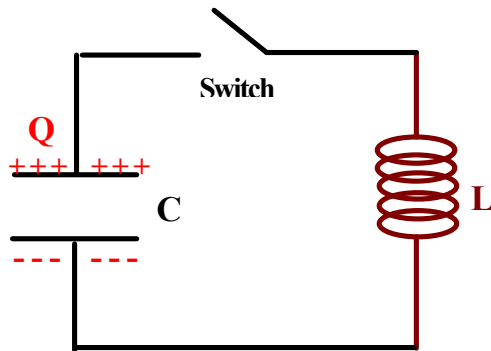
$$U_B = \frac{1}{2} LI^2$$

$$\Phi_B = LI \quad L = \Phi_B / I$$

$$\epsilon_L = -L \frac{dI}{dt}$$

$$u_B = \frac{1}{2\mu_0} B^2 \quad (\text{B-field energy density})$$

An LC Circuit



At $t = 0$ the switch is closed and a capacitor with initial charge Q_0 is connected in series across an inductor (assume there is no resistance). The initial conditions are $Q(0) = Q_0$ and $I(0) = 0$. Moving around the circuit in the direction of the current flow yields

$$\frac{Q}{C} - L \frac{dI}{dt} = 0.$$

Since I is flowing out of the capacitor, $I = -dQ/dt$, so that

$$\frac{d^2 Q}{dt^2} + \frac{1}{LC} Q = 0.$$

This differential equation for $Q(t)$ is the **SHM differential equation** we studied earlier with $\omega = \sqrt{1/LC}$ and solution

$$Q(t) = A \cos \omega t + B \sin \omega t.$$

The current is thus,

$$I(t) = -\frac{dQ}{dt} = A\omega \sin \omega t - B\omega \cos \omega t.$$

Applying the initial conditions yields

$$Q(t) = Q_0 \cos \omega t$$

$$I(t) = Q_0 \omega \sin \omega t$$

Thus, $Q(t)$ and $I(t)$ **oscillate with SHM with angular frequency**

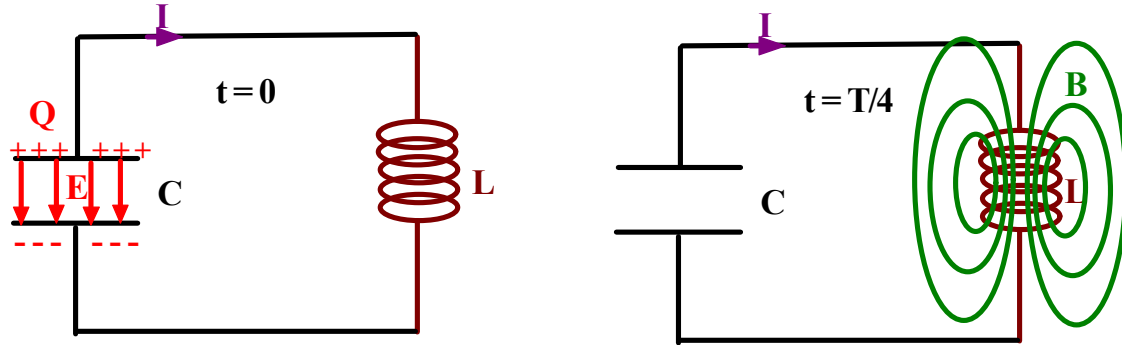
$\omega = \sqrt{1/LC}$. The stored energy oscillates between **electric** and **magnetic** according to

$$U_E(t) = \frac{Q^2(t)}{2C} = \frac{Q_0^2}{2C} \cos^2 \omega t$$

$$U_B(t) = \frac{1}{2} L I^2(t) = \frac{1}{2} L Q_0^2 \omega^2 \sin^2 \omega t$$

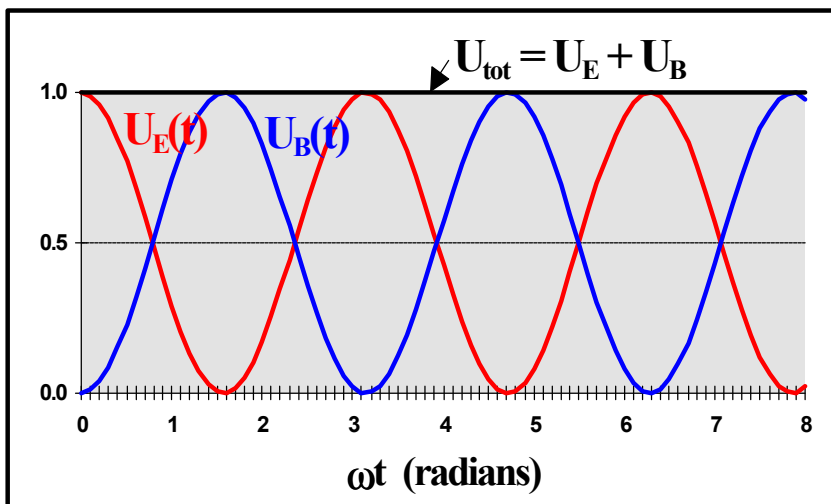
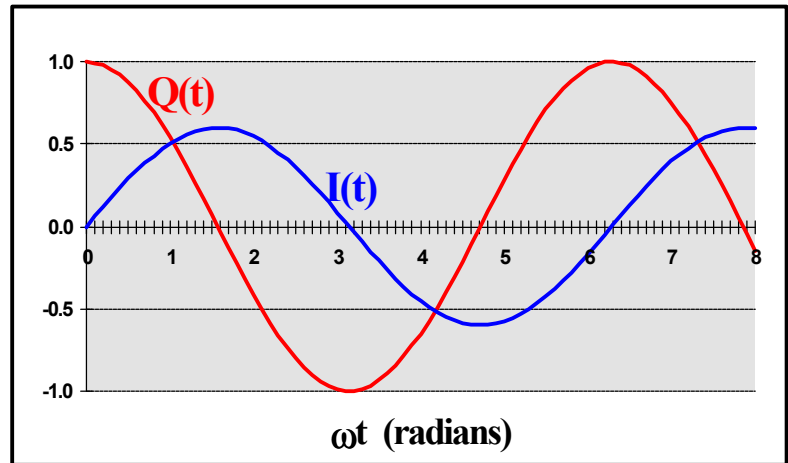
Energy is conserved since $U_{\text{tot}}(t) = U_E(t) + U_B(t) = Q_0^2/2C$ is **constant**.

LC Oscillations



$$Q(t) = Q_0 \cos \omega t$$

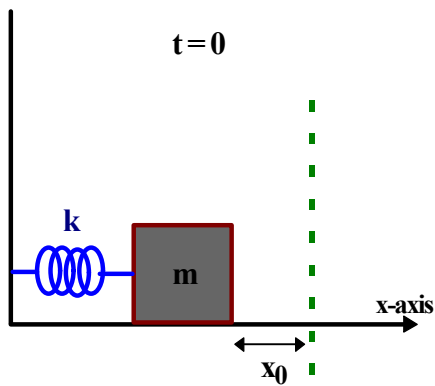
$$I(t) = Q_0 \omega \sin \omega t$$



$$U_E(t) = \frac{Q_0^2}{2C} \cos^2 \omega t$$

$$U_B(t) = \frac{Q_0^2}{2C} \sin^2 \omega t$$

Mechanical Analogy



At $t = 0$:

$$E = \frac{1}{2} k x_0^2$$

$$v = 0$$

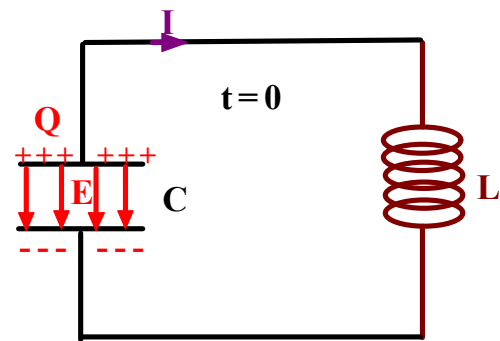
At Later t :

$$v = \frac{dx}{dt}$$

$$x(t) = x_0 \cos \omega t$$

$$\omega = \sqrt{\frac{k}{m}}$$

$$E = \frac{1}{2} m v^2 + \frac{1}{2} k x^2$$



At $t = 0$:

$$U = \frac{1}{2C} Q^2$$

$$I = 0$$

At Later t :

$$I = -\frac{dQ}{dt}$$

$$Q(t) = Q_0 \cos \omega t$$

$$\omega = \sqrt{\frac{1}{LC}}$$

$$E = \frac{1}{2} L I^2 + \frac{1}{2C} Q^2$$

Constant

Correspondence:

$$x(t) \leftrightarrow Q(t)$$

$$v(t) \leftrightarrow I(t)$$

$$m \leftrightarrow L$$

$$k \leftrightarrow 1/C$$

Another Differential Equation

Consider the **2nd order** differential equation

$$\frac{d^2 x(t)}{dt^2} + D \frac{dx(t)}{dt} + Cx(t) = 0,$$

where **C** and **D** are constants. We solve this equation by turning it into an

algebraic equation by looking for a solution of the form $x(t) = Ae^{at}$.

Substituting this into the differential equation yields,

$$a^2 + Da + C = 0 \quad \text{or} \quad \boxed{a = -\frac{D}{2} \pm \sqrt{\left(\frac{D}{2}\right)^2 - C}}.$$

Case I ($C > (D/2)^2$, damped oscillations):

For $C > (D/2)^2$, $a = -D/2 \pm i\sqrt{C - (D/2)^2} = -D/2 \pm i\omega'$, where

$\omega' = \sqrt{C - (D/2)^2}$, and the most general solution has the form:

$$\begin{aligned} x(t) &= e^{-Dt/2} (Ae^{i\omega't} + Be^{-i\omega't}) \\ x(t) &= e^{-Dt/2} (A \cos(\omega't) + B \sin(\omega't)) \\ x(t) &= Ae^{-Dt/2} \sin(\omega't + \phi) \\ x(t) &= Ae^{-Dt/2} \cos(\omega't + \phi) \end{aligned}$$

where **A**, **B**, and ϕ are arbitrary constants.

Case II ($C < (D/2)^2$, over damped):

For $C < (D/2)^2$, $a = -D/2 \pm \sqrt{(D/2)^2 - C} = -D/2 \pm \gamma$, where

$\gamma = \sqrt{(D/2)^2 - C}$. In this case,

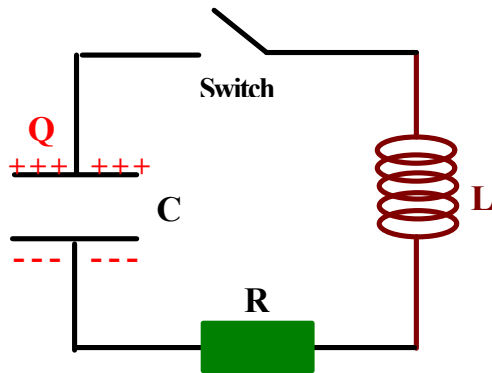
$$x(t) = e^{-Dt/2} (Ae^{\gamma t} + Be^{-\gamma t}).$$

Case III ($C = (D/2)^2$, critically damped):

For $C = (D/2)^2$, $a = -D/2$, and

$$x(t) = Ae^{-Dt/2}.$$

An LRC Circuit



At $t = 0$ the switch is closed and a capacitor with initial charge Q_0 is connected in series across an inductor and a resistor. The initial conditions are $Q(0) = Q_0$ and $I(0) = 0$. Moving around the circuit in the direction of the current flow yields

$$\frac{Q}{C} - L \frac{dI}{dt} - IR = 0.$$

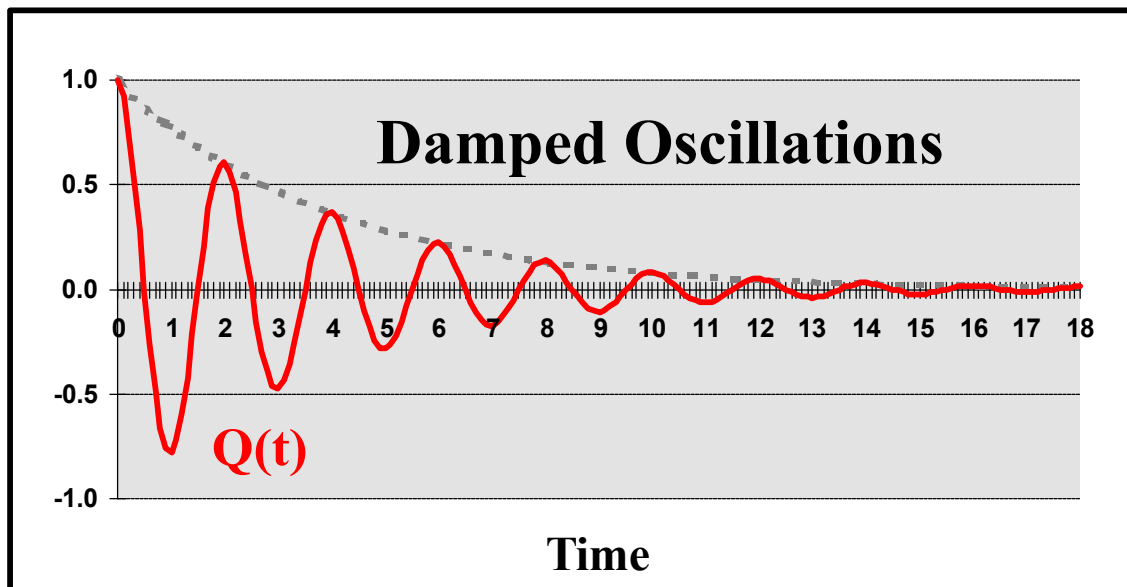
Since I is flowing out of the capacitor, $I = -dQ/dt$, so that

$$\frac{d^2Q}{dt^2} + \frac{R}{L} \frac{dQ}{dt} + \frac{1}{LC} Q = 0.$$

This differential equation for $Q(t)$ is the differential equation we studied earlier. If we take the case where $R^2 < 4L/C$ (**damped oscillations**) then

$$Q(t) = Q_0 e^{-Rt/2L} \cos \omega' t,$$

with $\omega' = \sqrt{\omega^2 - (R/2L)^2}$ and $\omega = \sqrt{1/LC}$.

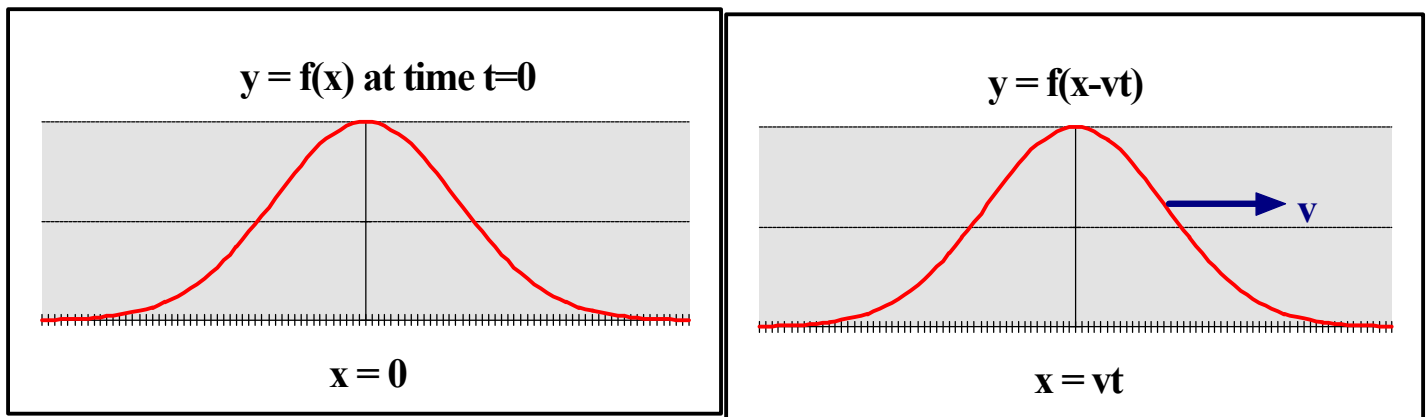


Traveling Waves

A **"wave"** is a traveling disturbance that transports energy but not matter.

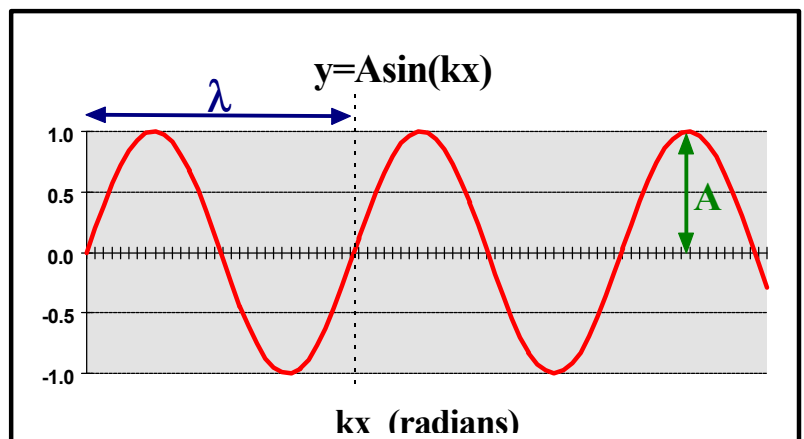
Constructing Traveling Waves:

To construct a wave with shape $y = f(x)$ at time $t = 0$ traveling to the **right** with speed v simply make the replacement $x \rightarrow x - vt$.



Traveling Harmonic Waves:

Harmonic waves have the form $y = A \sin(kx)$ or $y = A \cos(kx)$ at time $t = 0$, where k is the **"wave number"** ($k = 2\pi/\lambda$ where λ is the **"wave length"**) and A is the **"amplitude"**. To construct an harmonic wave traveling to the right with speed v , replace x by $x-vt$ as follows:



$y = A \sin(k(x-vt)) = A \sin(kx - \omega t)$ where $\omega = kv$ ($v = \omega/k$). The period of the oscillation, $T = 2\pi/\omega = 1/f$, where f is the **linear frequency** (measured in **Hertz** where $1\text{Hz} = 1/\text{sec}$) and ω is the **angular frequency** ($\omega = 2\pi f$). The speed of propagation is given by $v = \omega/k = \lambda f$.

$y = y(x,t) = A \sin(kx - \omega t)$ right moving harmonic wave

$y = y(x,t) = A \sin(kx + \omega t)$ left moving harmonic wave

The Wave Equation

$$\boxed{\frac{\partial^2 y(x,t)}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2 y(x,t)}{\partial t^2} = 0}$$

Whenever analysis of a system results in an equation of the form given above then we know that the system supports traveling waves propagating at speed v .

General Proof:

If $y = y(x,t) = f(x-vt)$ then

$$\begin{aligned} \frac{\partial y}{\partial x} &= f' & \frac{\partial^2 y}{\partial x^2} &= f'' \\ \frac{\partial y}{\partial t} &= -vf' & \frac{\partial^2 y}{\partial t^2} &= v^2 f'' \end{aligned}$$

and

$$\frac{\partial^2 y(x,t)}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2 y(x,t)}{\partial t^2} = f'' - f'' = 0.$$

Proof for Harmonic Wave:

If $y = y(x,t) = A \sin(kx - \omega t)$ then

$$\frac{\partial^2 y}{\partial x^2} = -k^2 A \sin(kx - \omega t) \quad \frac{\partial^2 y}{\partial t^2} = -\omega^2 A \sin(kx - \omega t)$$

and

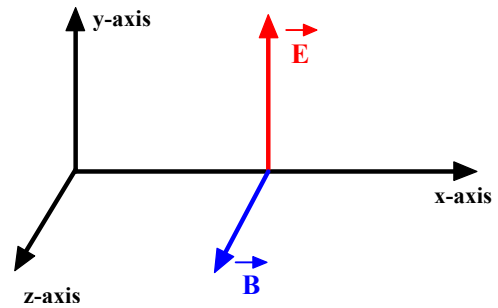
$$\frac{\partial^2 y(x,t)}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2 y(x,t)}{\partial t^2} = \left(-k^2 + \frac{\omega^2}{v^2} \right) A \sin(kx - \omega t) = 0,$$

since $\omega = kv$.

Light Propagating in Empty Space

Since there are no charges and no current in empty space, Faraday's Law and Ampere's Law take the form

$$\oint \vec{E} \cdot d\vec{l} = -\frac{d\Phi_B}{dt} \quad \oint \vec{B} \cdot d\vec{l} = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}.$$



Look for a solution of the form

$$\vec{E}(x, t) = E_y(x, t) \hat{y}$$

$$\vec{B}(x, t) = B_z(x, t) \hat{z}$$

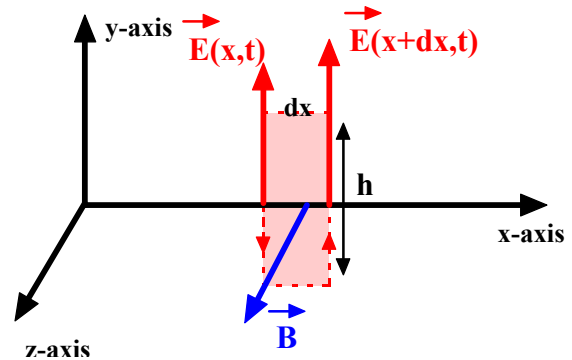
Faraday's Law:

Computing the left and right hand side of Faraday's Law using a rectangle (in the xy-plane) with width dx and height h (counterclockwise) gives

$$E_y(x + dx, t)h - E_y(x, t)h = -\frac{\partial B_z}{\partial t} h dx$$

or

$$\boxed{\frac{\partial E_y}{\partial x} = -\frac{\partial B_z}{\partial t}}$$



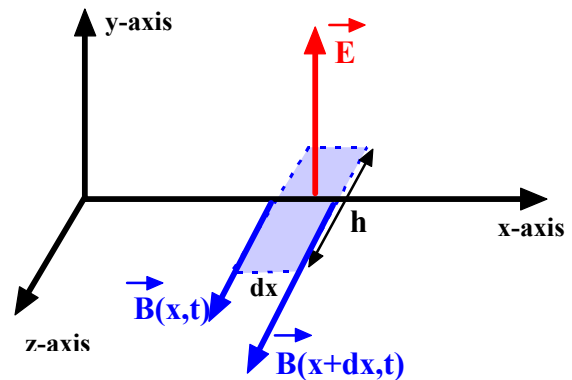
Ampere's Law:

Computing the left and right hand side of Ampere's Law using a rectangle (in the xz-plane) with width dx and height h (counterclockwise) gives

$$B_z(x, t)h - B_z(x + dx, t)h = \mu_0 \epsilon_0 \frac{\partial E_y}{\partial t} h dx$$

or

$$\boxed{-\frac{\partial B_z}{\partial x} = \mu_0 \epsilon_0 \frac{\partial E_y}{\partial t}}$$



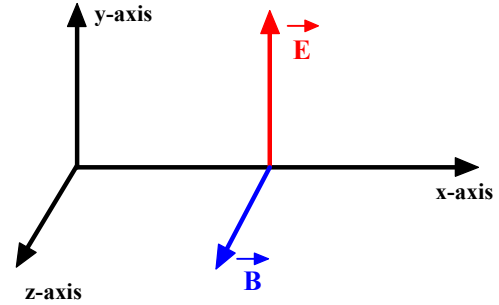
Electromagnetic Plane Waves (1)

We have the following two **differential equations** for $E_y(x,t)$ and $B_z(x,t)$:

$$\frac{\partial B_z}{\partial t} = -\frac{\partial E_y}{\partial x} \quad (1)$$

and

$$\frac{\partial E_y}{\partial t} = -\frac{1}{\mu_0 \epsilon_0} \frac{\partial B_z}{\partial x} \quad (2)$$



Taking the time derivative of (2) and using (1) gives

$$\frac{\partial^2 E_y}{\partial t^2} = -\frac{1}{\mu_0 \epsilon_0} \frac{\partial}{\partial t} \left(\frac{\partial B_z}{\partial x} \right) = -\frac{1}{\mu_0 \epsilon_0} \frac{\partial}{\partial x} \left(\frac{\partial B_z}{\partial t} \right) = \frac{1}{\mu_0 \epsilon_0} \frac{\partial^2 E_y}{\partial x^2}$$

which implies

$$\frac{\partial^2 E_y}{\partial x^2} - \mu_0 \epsilon_0 \frac{\partial^2 E_y}{\partial t^2} = 0.$$

Thus $E_y(x,t)$ satisfies the wave equation with speed $v = 1/\sqrt{\epsilon_0 \mu_0}$ and has a solution in the form of traveling waves as follows:

$$\mathbf{E}_y(x,t) = \mathbf{E}_0 \sin(kx - \omega t),$$

where \mathbf{E}_0 is the **amplitude of the electric field oscillations** and where the wave has a **unique speed**

$$v = c = \frac{\omega}{k} = \lambda f = \frac{1}{\sqrt{\epsilon_0 \mu_0}} = 2.99792 \times 10^8 \text{ m/s (speed of light).}$$

From (1) we see that

$$\frac{\partial B_z}{\partial t} = -\frac{\partial E_y}{\partial x} = -E_0 k \cos(kx - \omega t),$$

which has a solution given by

$$B_z(x,t) = E_0 \frac{k}{\omega} \sin(kx - \omega t) = \frac{E_0}{c} \sin(kx - \omega t),$$

so that

$$\mathbf{B}_z(x,t) = \mathbf{B}_0 \sin(kx - \omega t),$$

where $\mathbf{B}_0 = \mathbf{E}_0/c$ is the **amplitude of the magnetic field oscillations**.

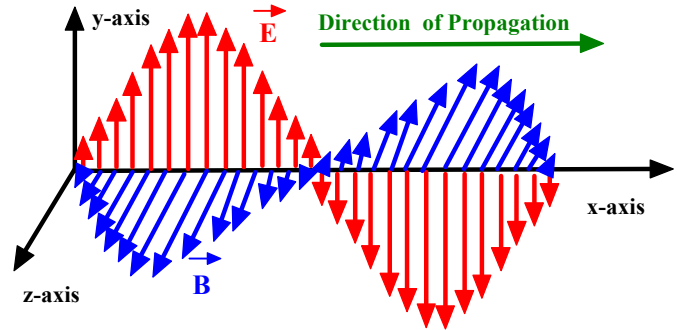
Electromagnetic Plane Waves (2)

The **plane harmonic wave solution** for light with frequency f and wavelength λ and speed $c = f\lambda$ is given by

$$\vec{E}(x, t) = E_0 \sin(kx - \omega t) \hat{y}$$

$$\vec{B}(x, t) = B_0 \sin(kx - \omega t) \hat{z}$$

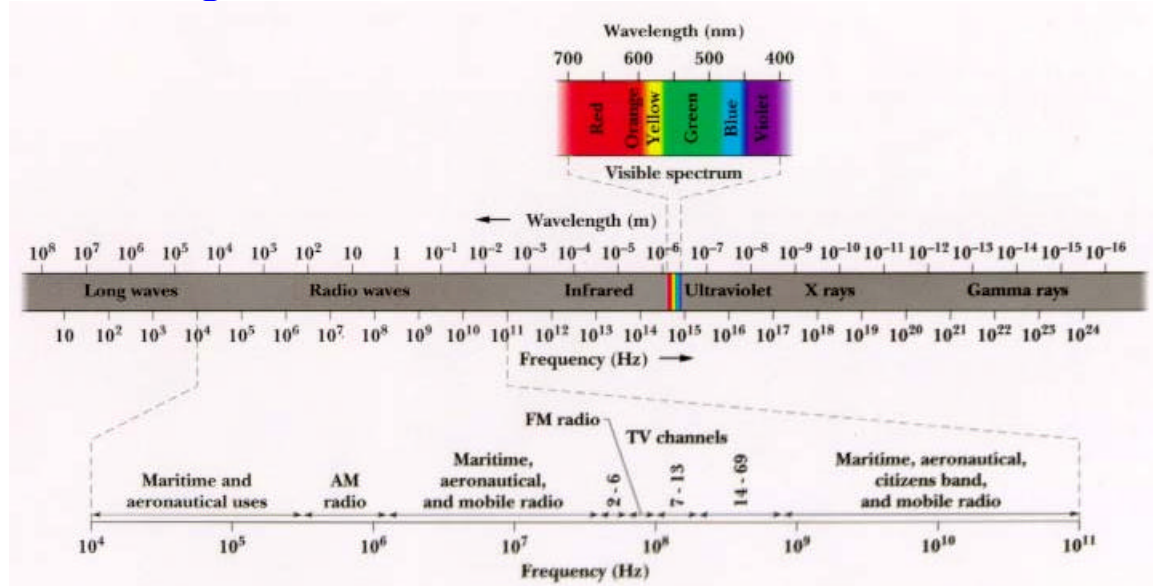
where $k = 2\pi/\lambda$, $\omega = 2\pi f$, and $E_0 = cB_0$.



Properties of the Electromagnetic Plane Wave:

- **Wave travels at speed c ($c = 1/\sqrt{\mu_0\epsilon_0}$).**
- **E and B are perpendicular ($\vec{E} \cdot \vec{B} = 0$).**
- **The wave travels in the direction of \vec{E} \vec{B} .**
- **At any point and time $E = cB$.**

Electromagnetic Radiation:



Energy Transport - Poynting Vector

Electric and Magnetic Energy Density:

For an **electromagnetic plane wave**

$$\mathbf{E}_y(\mathbf{x},t) = \mathbf{E}_0 \sin(kx - \omega t),$$

$$\mathbf{B}_z(\mathbf{x},t) = \mathbf{B}_0 \sin(kx - \omega t),$$

where $\mathbf{B}_0 = \mathbf{E}_0/c$. The **electric energy density** is given by

$$u_E = \frac{1}{2} \epsilon_0 E^2 = \frac{1}{2} \epsilon_0 E_0^2 \sin^2(kx - \omega t) \text{ and the magnetic energy density is}$$

$$u_B = \frac{1}{2\mu_0} B^2 = \frac{1}{2\mu_0 c^2} E^2 = \frac{1}{2} \epsilon_0 E^2 = u_E,$$

where I used $\mathbf{E} = c\mathbf{B}$. Thus, **for light the electric and magnetic field energy densities are equal** and the **total energy density** is

$$u_{tot} = u_E + u_B = \epsilon_0 E^2 = \frac{1}{\mu_0} B^2 = \epsilon_0 E_0^2 \sin^2(kx - \omega t).$$

Poynting Vector ($\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B}$):

The **direction** of the **Poynting Vector** is the **direction of energy flow** and the **magnitude**

$$S = \frac{1}{\mu_0} EB = \frac{E^2}{\mu_0 c} = \frac{1}{A} \frac{dU}{dt}$$

is the **energy per unit time per unit area** (units of **Watts/m²**).

Proof:

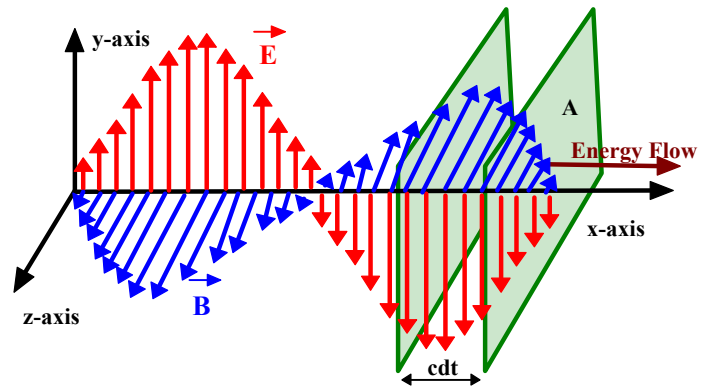
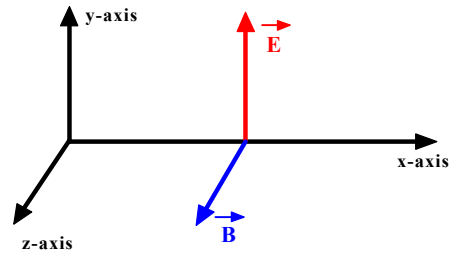
$$dU_{tot} = u_{tot} V = \epsilon_0 E^2 A c dt \text{ so}$$

$$S = \frac{1}{A} \frac{dU}{dt} = \epsilon_0 c E^2 = \frac{E^2}{\mu_0 c} = \frac{E_0^2}{\mu_0 c} \sin^2(kx - \omega t).$$

Intensity of the Radiation (Watts/m²):

The intensity, **I**, is the **average of S** as follows:

$$I = \bar{S} = \frac{1}{A} \frac{d\bar{U}}{dt} = \frac{E_0^2}{\mu_0 c} \langle \sin^2(kx - \omega t) \rangle = \frac{E_0^2}{2\mu_0 c}.$$



Momentum Transport - Radiation Pressure

Relativistic Energy and Momentum:

$$E^2 = (cp)^2 + (m_0c^2)^2$$

↑ energy ↑ momentum ↑ rest mass

$$E = cp \quad (\text{for light})$$

For light $m_0 = 0$ and

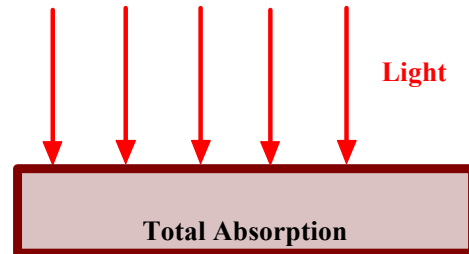
For light the **average momentum per unit time per unit area** is equal to the intensity of the light, I , divided by speed of light, c , as follows:

$$\frac{1}{A} \frac{d\bar{p}}{dt} = \frac{1}{c} \frac{1}{A} \frac{d\bar{U}}{dt} = \frac{1}{c} I.$$

Total Absorption:

$$\bar{F} = \frac{d\bar{p}}{dt} = \frac{1}{c} \frac{d\bar{U}}{dt} = \frac{1}{c} IA$$

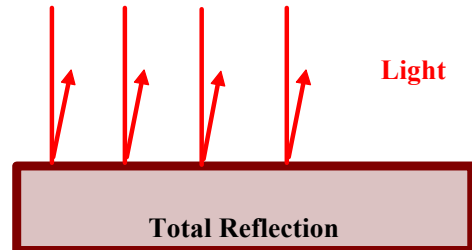
$$P = \frac{\bar{F}}{A} = \frac{1}{c} I \quad (\text{radiation pressure})$$



Total Reflection:

$$\bar{F} = \frac{d\bar{p}}{dt} = \frac{2}{c} \frac{d\bar{U}}{dt} = \frac{2}{c} IA$$

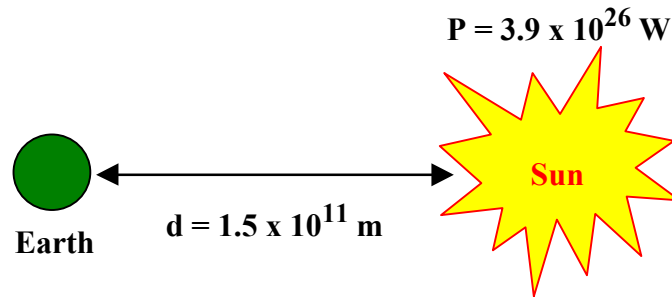
$$P = \frac{\bar{F}}{A} = \frac{2}{c} I \quad (\text{radiation pressure})$$



The Radiation Power of the Sun

Problem:

The **radiation power** of the sun is $3.9 \times 10^{26} \text{ W}$ and the distance from the Earth to the sun is $1.5 \times 10^{11} \text{ m}$.



- (a) What is the **intensity** of the electromagnetic radiation from the sun at the surface of the Earth (outside the atmosphere)? (**answer: 1.4 kW/m^2**)
- (b) What is the **maximum value of the electric field** in the light coming from the sun? (**answer: $1,020 \text{ V/m}$**)
- (c) What is the **maximum energy density of the electric field** in the light coming from the sun? (**answer: $4.6 \times 10^{-6} \text{ J/m}^3$**)
- (d) What is the **maximum value of the magnetic field** in the light coming from the sun? (**answer: 3.4 T**)
- (e) What is the **maximum energy density of the magnetic field** in the light coming from the sun? (**answer: $4.6 \times 10^{-6} \text{ J/m}^3$**)
- (f) Assuming complete absorption what is the **radiation pressure on the Earth** from the light coming from the sun? (**answer: $4.7 \times 10^{-6} \text{ N/m}^2$**)
- (g) Assuming complete absorption what is **the radiation force on the Earth** from the light coming from the sun? The radius of the Earth is about $6.4 \times 10^6 \text{ m}$. (**answer: $6 \times 10^8 \text{ N}$**)
- (h) What is the **gravitational force** on the Earth due to the sun. The mass of the Earth and the sun are $5.98 \times 10^{24} \text{ kg}$ and $1.99 \times 10^{30} \text{ kg}$, respectively, and $G = 6.67 \times 10^{-11} \text{ Nm}^2/\text{kg}^2$. (**answer: $3.5 \times 10^{22} \text{ N}$**)